

March 12-16, 2018

ICERM Workshop on “Fast Algorithms for Generating Static and Dynamically Changing Point Configurations”

Improving Particle Methods

Robert Krasny (Michigan)

collaborators

Weihua Geng (Southern Methodist University)

Peter Bosler (Sandia National Laboratories)

Ling Xu (Michigan)

support

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MICDE (Michigan Institute for Computational Discovery and Engineering)

outline

1. protein-solvent electrostatics
2. incompressible fluid flow

1. protein-solvent electrostatics

atomic charges

Van der Waals radius

water molecules

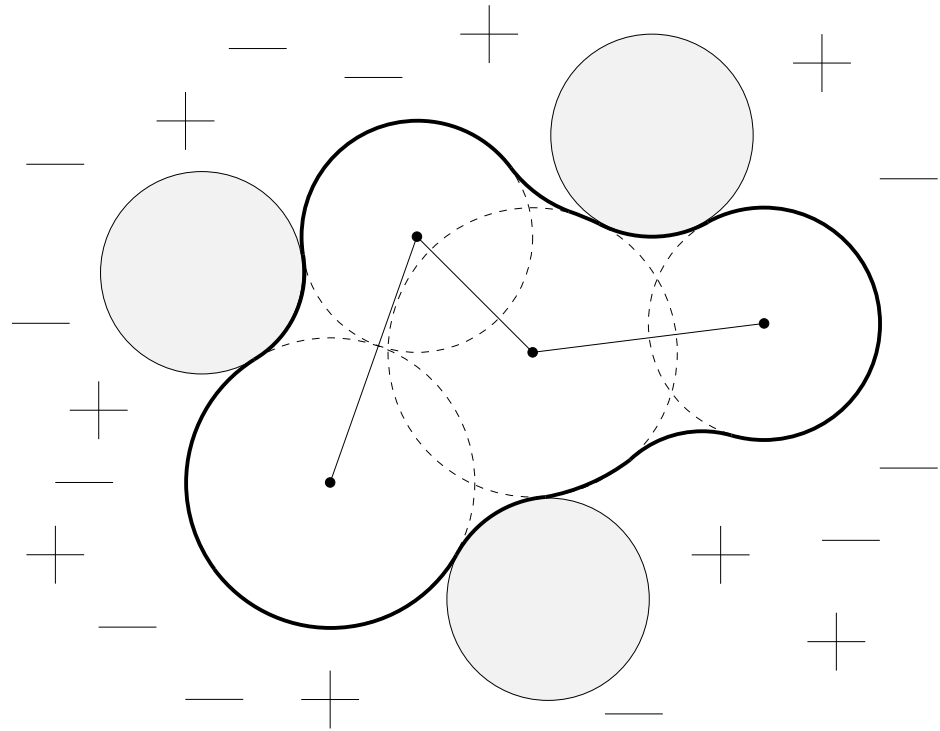
molecular surface

dissolved salt ions

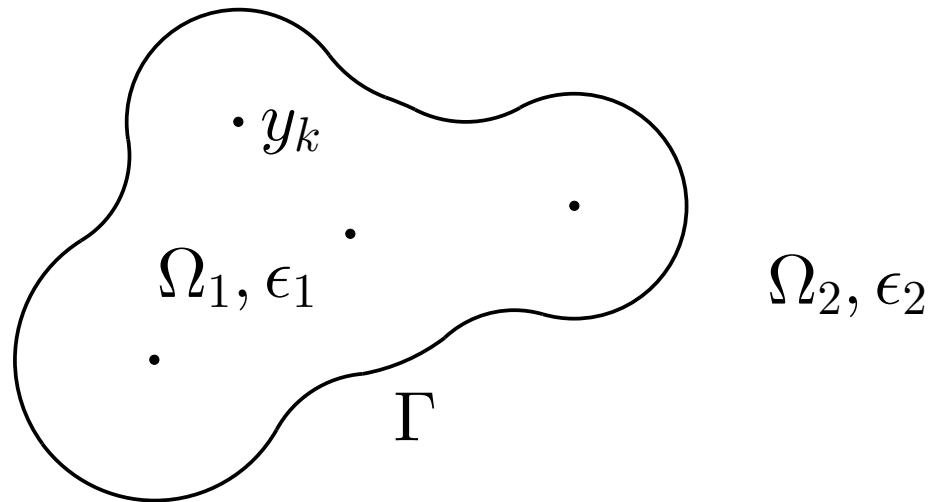
goal

$\phi(x)$: electrostatic potential

E_{sol} : electrostatic solvation energy



Poisson-Boltzmann implicit solvent model



$$\Omega_1 : \text{protein domain} : \epsilon_1 \nabla^2 \phi(x) = - \sum_{k=1}^{N_c} q_k \delta(x - y_k)$$

$$\Omega_2 : \text{solvent domain} : \epsilon_2 \nabla^2 \phi(x) - \bar{\kappa}^2 \phi(x) = 0 : \text{PB equation}$$

$$\Gamma : \text{molecular surface} : \phi_1 = \phi_2, \quad \epsilon_1 \frac{\partial \phi_1}{\partial n} = \epsilon_2 \frac{\partial \phi_2}{\partial n}$$

$$\text{far-field boundary condition} : \lim_{|x| \rightarrow \infty} \phi(x) = 0$$

boundary integral formulation Juffer et al. (1991)

$$G_0(x, y) = \frac{1}{4\pi|x - y|} : \text{Coulomb potential}$$

$$G_\kappa(x, y) = \frac{e^{-\kappa|x-y|}}{4\pi|x - y|} : \text{screened Coulomb potential} , \quad \kappa^2 = \bar{\kappa}^2/\epsilon_2$$

$$\frac{1}{2}\left(1 + \frac{\epsilon_2}{\epsilon_1}\right)\phi(x) = \int_{\Gamma} \left[K_1(x, y) \partial_n \phi(y) + K_2(x, y) \phi(y) \right] dS_y + S_1(x)$$

$$\frac{1}{2}\left(1 + \frac{\epsilon_1}{\epsilon_2}\right) \partial_n \phi(x) = \int_{\Gamma} \left[K_3(x, y) \partial_n \phi(y) + K_4(x, y) \phi(y) \right] dS_y + S_2(x)$$

$$K_1 = G_0 - G_\kappa , \quad K_2 = \frac{\epsilon_2}{\epsilon_1} \partial_{n_y} G_\kappa - \partial_{n_y} G_0$$

$$K_4 = \partial_{n_x n_y}^2 (G_\kappa - G_0) , \quad K_3 = \partial_{n_x} G_0 - \frac{\epsilon_1}{\epsilon_2} \partial_{n_x} G_\kappa$$

$$S_1(x) = \frac{1}{\epsilon_1} \sum_{k=1}^{N_c} q_k G_0(x, y_k) , \quad S_2(x) = \frac{1}{\epsilon_1} \sum_{k=1}^{N_c} q_k \partial_{n_x} G_0(x, y_k)$$

discretization

triangulation + centroid collocation

x_i : centroid , A_i : area , $i = 1 : N$

$$\frac{1}{2} \left(1 + \frac{\epsilon_2}{\epsilon_1} \right) \phi_i = \sum_{\substack{j=1 \\ j \neq i}}^N \left[K_1(x_i, x_j) \partial_n \phi_j + K_2(x_i, x_j) \phi_j \right] A_j + S_1(x_i)$$

$$\frac{1}{2} \left(1 + \frac{\epsilon_1}{\epsilon_2} \right) \partial_n \phi_i = \sum_{\substack{j=1 \\ j \neq i}}^N \left[K_3(x_i, x_j) \partial_n \phi_j + K_4(x_i, x_j) \phi_j \right] A_j + S_2(x_i)$$

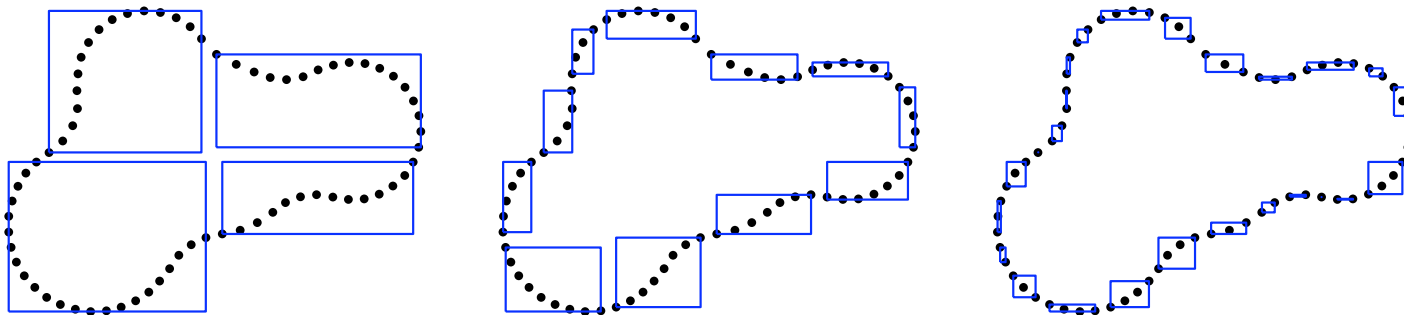
$\Rightarrow$ linear system : $Ax = b$

- GMRES
- matrix-vector product at each step
- fast summation $\rightarrow$ treecode

treecode : higher-order version of Barnes-Hut (1986)

$$\phi_i = \sum_{\substack{j=1 \\ j \neq i}}^N q_j G(x_i, x_j) , \quad i = 1, \dots, N : \text{particle-particle} , O(N^2)$$
$$\approx \sum_c \sum_{k=0}^p a_k(x_i, x_c) M^k(c) : \text{particle-cluster} , O(N \log N)$$

- Cartesian coordinates
- recurrence relations for $a_k(x_i, x_c) = \frac{1}{k!} D_y^k G(x, x_c)$
- geometrically adaptive

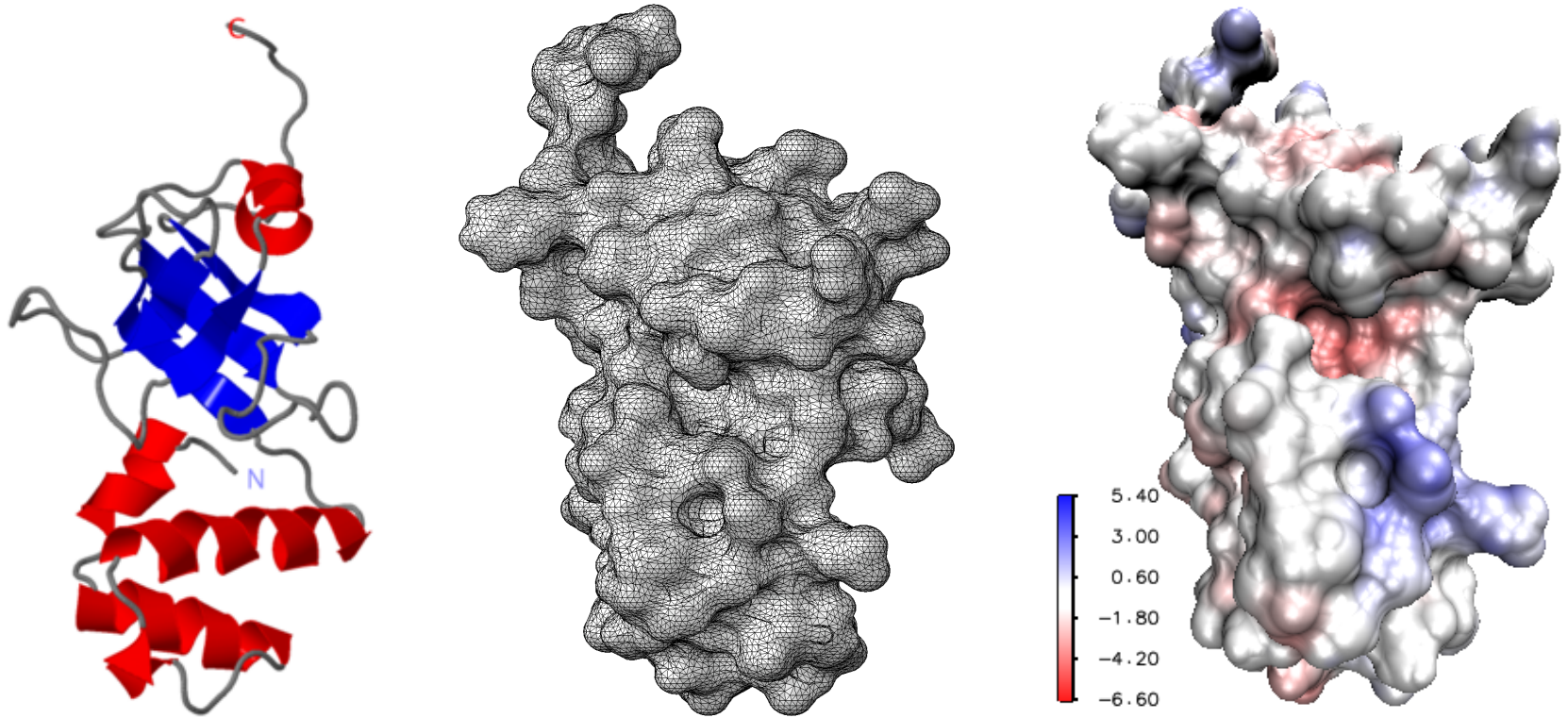


- low memory usage, good parallel efficiency

example : protein 1A63

RNA binding domain of E. coli rho factor

2069 atoms , triangulation by MSMS with $N = 132196$



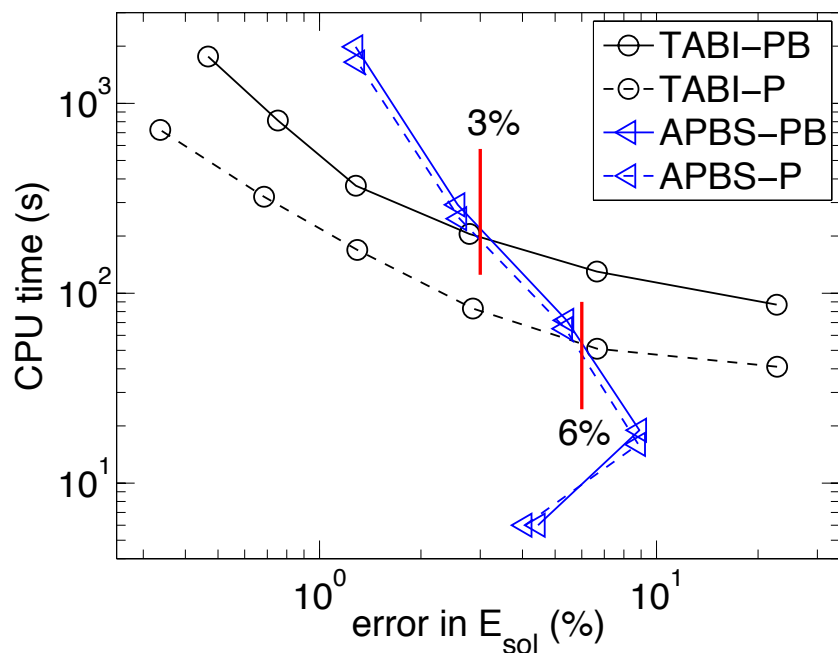
example : protein 1A63

compare 2 codes

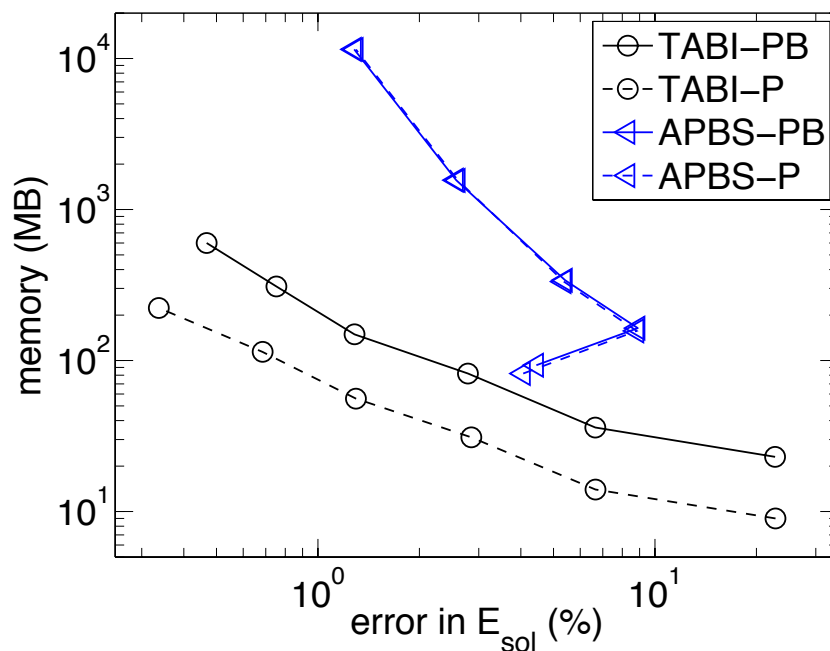
TABI : $N = 20\text{K} \rightarrow 500\text{K}$

APBS : grid = $65 \times 33^2 \rightarrow 513 \times 321^2$, $h_{\text{max}} = 1.63 \text{ \AA} \rightarrow 0.13 \text{ \AA}$

CPU time / error



memory / error



example : protein 1A63

TABI parallel performance , $N = 132196$, error in $E_{\text{sol}} \approx 1.3\%$

Poisson-Boltzmann

# processors	run time (s)	speedup	parallel efficiency (%)
1	799.3	1.00	100.0
2	410.0	1.95	97.5
4	223.8	3.57	89.3
8	123.7	6.46	80.9

Poisson

# processors	run time (s)	speedup	parallel efficiency (%)
1	324.4	1.00	100.0
2	166.7	1.95	97.3
4	92.6	3.50	87.6
8	51.0	6.36	79.5



TOOLS FOR PROTEIN SCIENCE

Improvements to the APBS biomolecular solvation software suite

Elizabeth Jurrus,¹ Dave Engel,¹ Keith Star,¹ Kyle Monson,¹ Juan Brandi,¹ Lisa E. Felberg,² David H. Brookes,² Leighton Wilson,³ Jiahui Chen,⁴ Karina Liles,¹ Minju Chun,¹ Peter Li,¹ David W. Gohara,⁵ Todd Dolinsky,⁶ Robert Konecny,⁷ David R. Koes ,⁸ Jens Erik Nielsen,⁹ Teresa Head-Gordon,² Weihua Geng,⁴ Robert Krasny,³ Guo-Wei Wei,¹⁰ Michael J. Holst,⁷ J. Andrew McCammon,⁷ and Nathan A. Baker ^{1,11*}

¹Pacific Northwest National Laboratory, Richland, Washington

²University of California, Berkeley, California

³University of Michigan, Ann Arbor, Michigan

⁴Southern Methodist University, Dallas, Texas

⁵St. Louis University, St. Louis, Missouri

⁶FoodLogiQ, Durham, North Carolina

⁷University of California San Diego, San Diego, California

⁸University of Pittsburgh, Pittsburgh, Pennsylvania

⁹Protein Engineering, Novozymes A/S, Copenhagen, Denmark

¹⁰Michigan State University, East Lansing, Michigan

¹¹Brown University, Providence, Rhode Island

comparison of software for molecular surface

MSMS

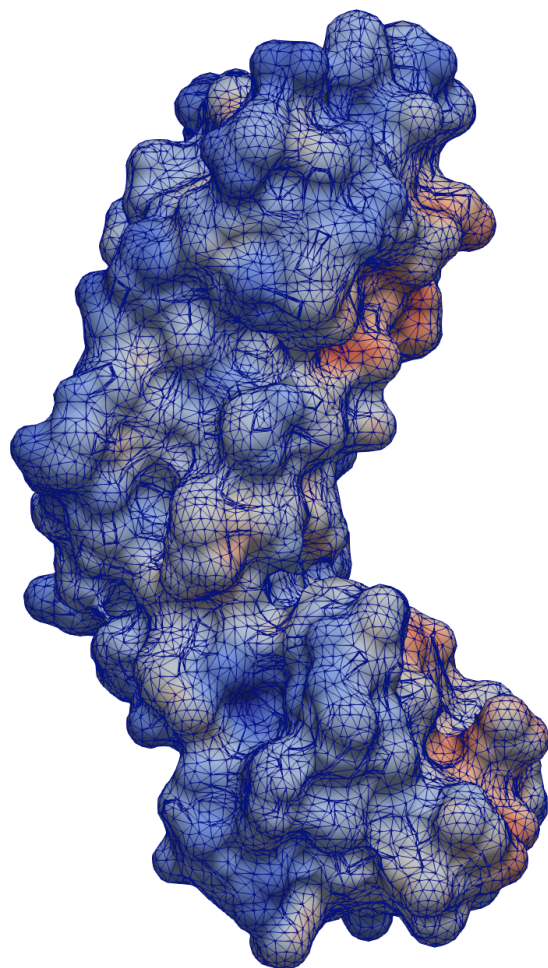
N = 29867, iter = 62

E_{sol} = -10443 kJ/mol

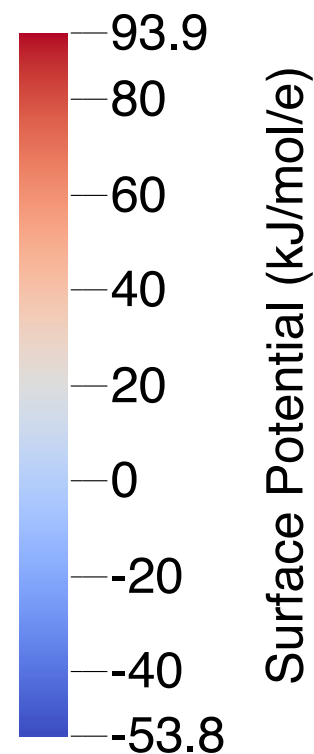
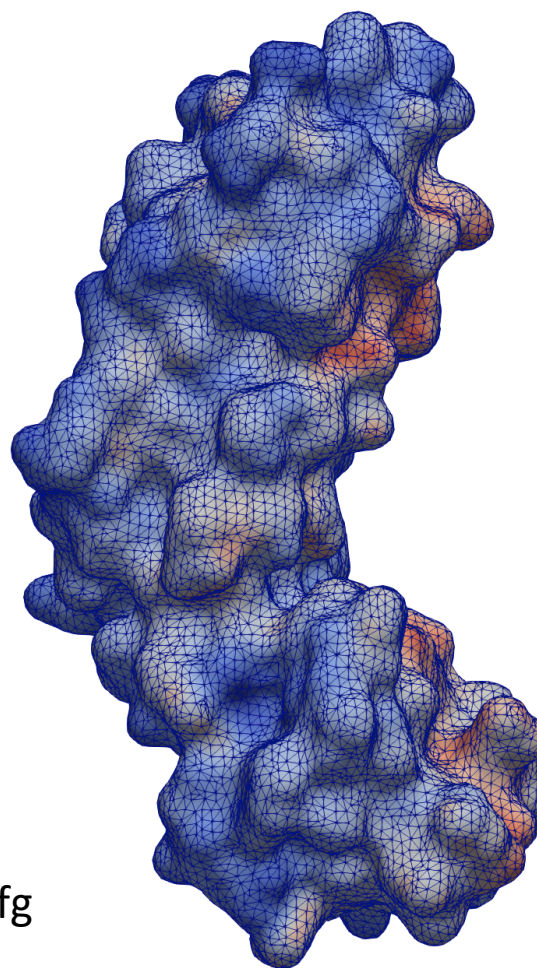
NanoShaper

N = 29224, iter = 9

E_{sol} = -10675 kJ/mol



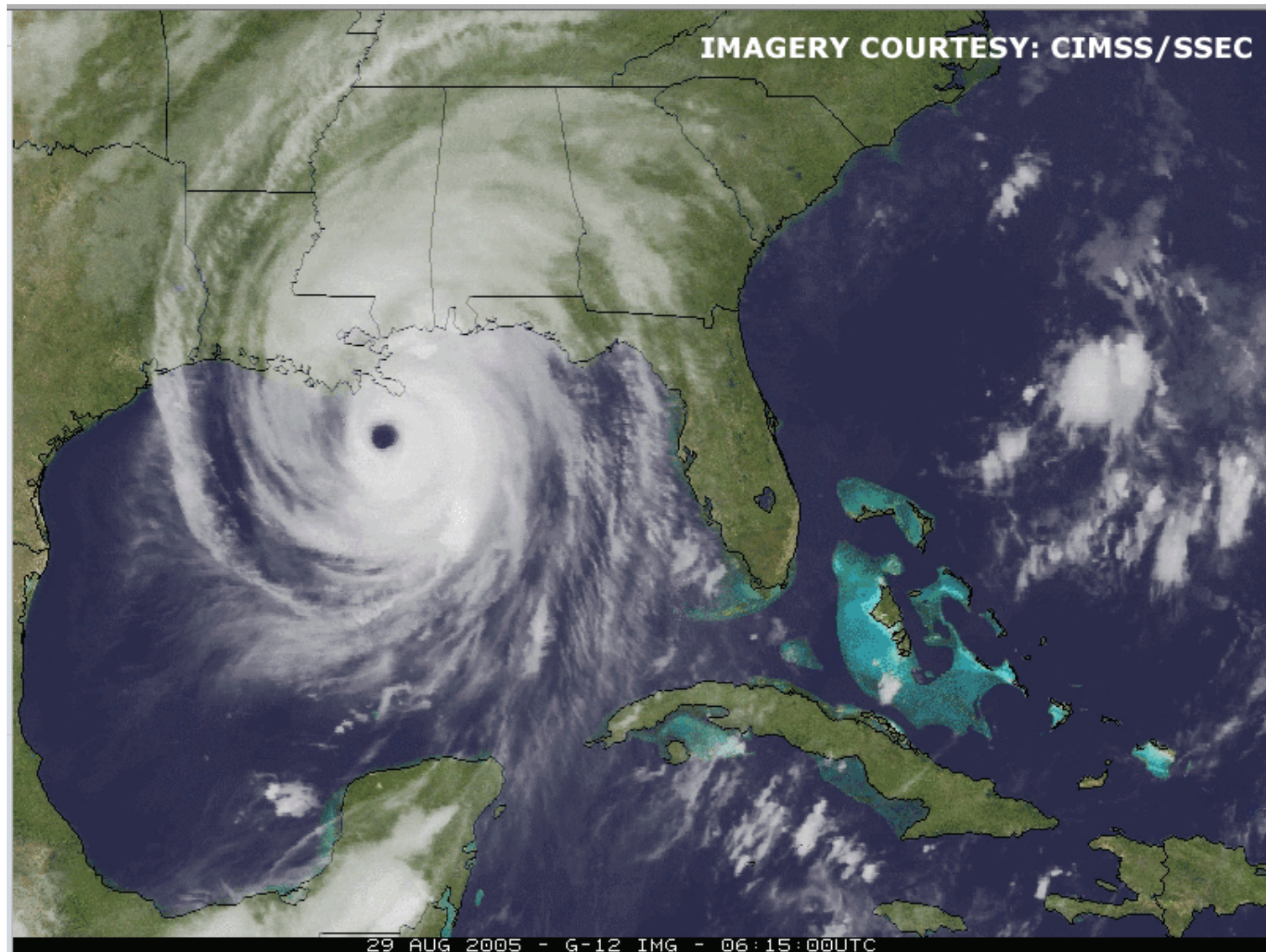
PDBID: 3dfg



outline

1. protein-solvent electrostatics
2. incompressible fluid flow
 - tracer transport on a sphere
 - vortex dynamics on a sphere
 - vortex dynamics in 2D

motivation : atmosphere



tracer transport on a sphere

given : $q_0(x)$, $u(x, t)$, $x \in S$

problem : determine $q(x, t)$ for $t > 0$

Eulerian form

$$\frac{\partial q}{\partial t}(x, t) + u \cdot \nabla q(x, t) = 0$$

Lagrangian form

flow map : $a \rightarrow x(a, t)$, $a \in S$

$$\frac{\partial}{\partial t} x(a, t) = u(x(a, t), t) , \quad x(a, 0) = a$$

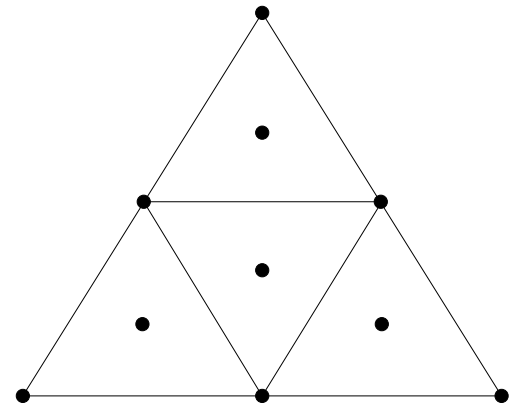
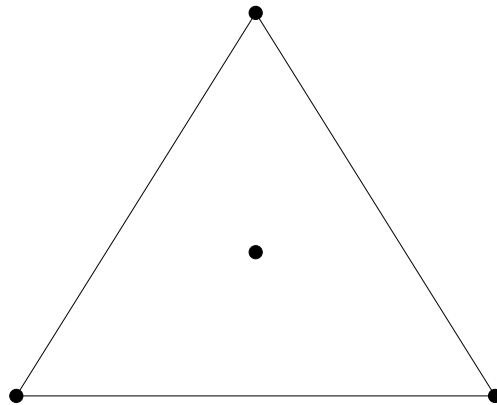
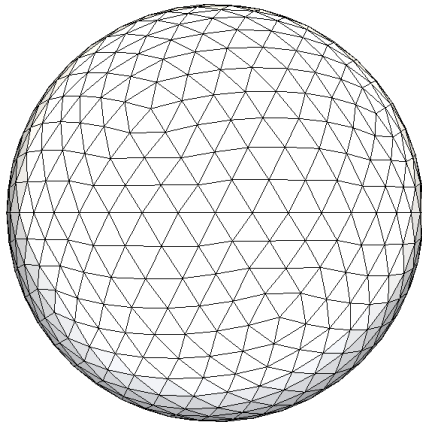
$$q(x(a, t), t) = q_0(a)$$

Lagrangian particle method

particles : $x_j(t) \approx x(a_j, t)$, $x_j(0) = a_j$, $j = 1 : M$

panels : $P_k(t) = x(P_k(0), t)$, $S = \cup_{k=1}^N P_k(t)$, $k = 1 : N$

icosahedral triangulation



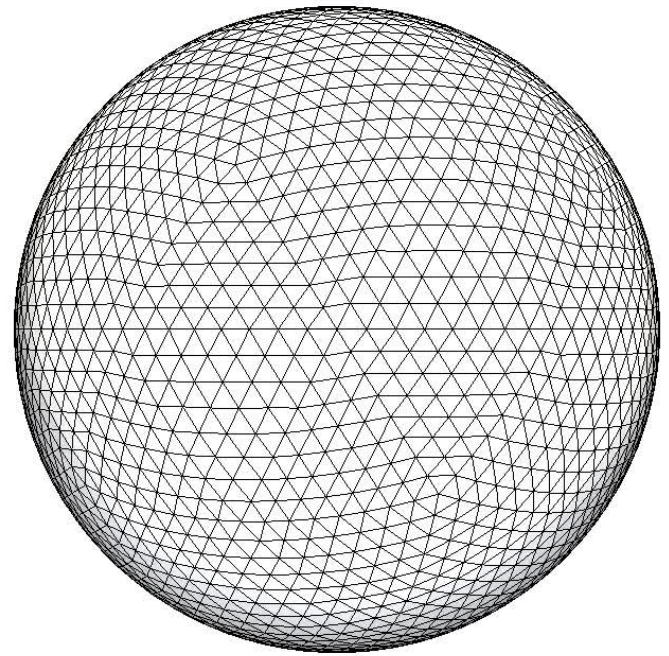
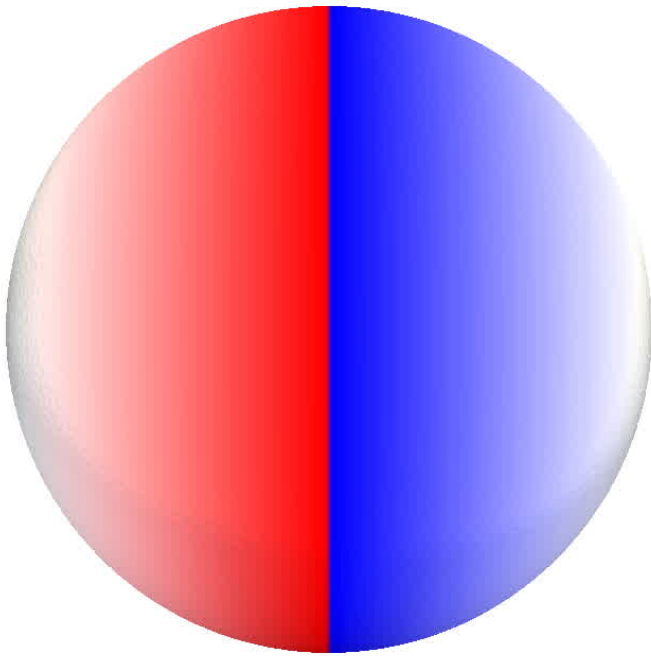
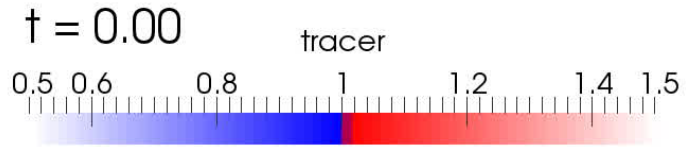
particle advection

$$\frac{d}{dt}x_j(t) = u(x_j, t) \text{ , } x_j(0) = a_j$$

$$q_j = q_0(a_j)$$

problem : particles become disordered for $t > 0$

test case 1 : moving vortices flow



movie

particle advection

$$\frac{d}{dt}x_j(t) = u(x_j, t) , \quad x_j(0) = a_j$$

$$q_j = q_0(a_j)$$

problem : particles become disordered for $t > 0$

solution : remeshing : $\{x_j^{old}\} \rightarrow \{x_j^{new}\} , \{q_j^{new}\} = ?$

1. direct remeshing

$$\{q_j^{new}\} = I(\{x_j^{new}\}; \{x_j^{old}\}, \{q_j^{old}\})$$

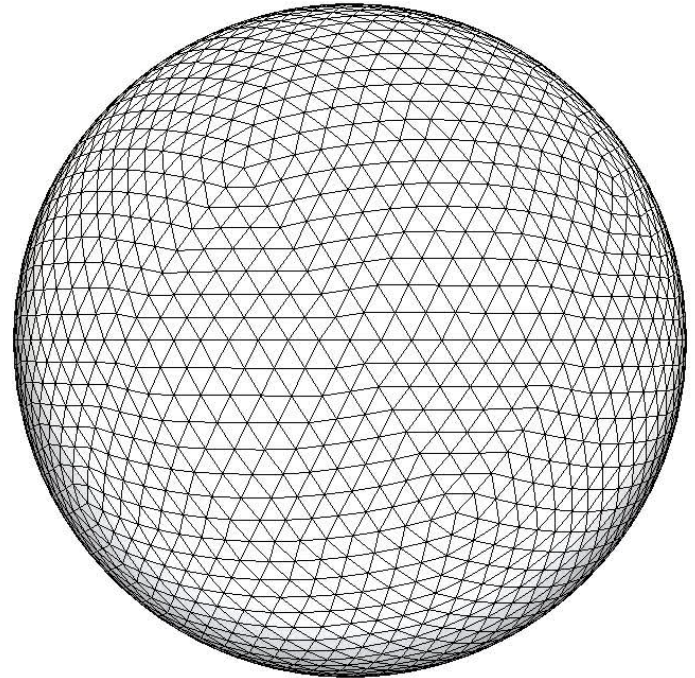
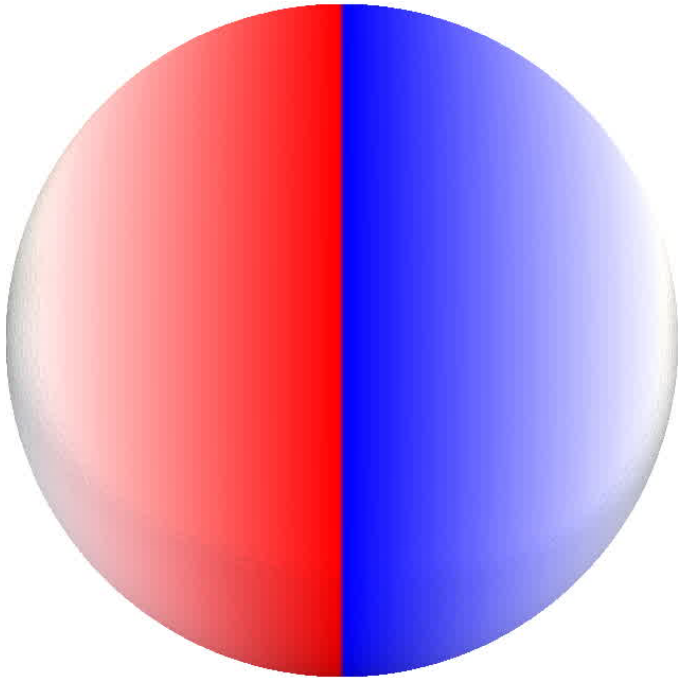
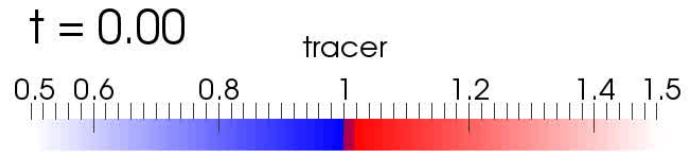
2. indirect remeshing

$$\{a_j^{new}\} = I(\{x_j^{new}\}; \{x_j^{old}\}, \{a_j^{old}\})$$

$$\{q_j^{new}\} = q_0(\{a_j^{new}\})$$

I : STRIPACK/SSRFPACK (Renka 1997)

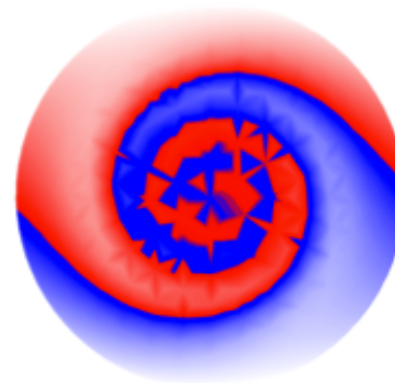
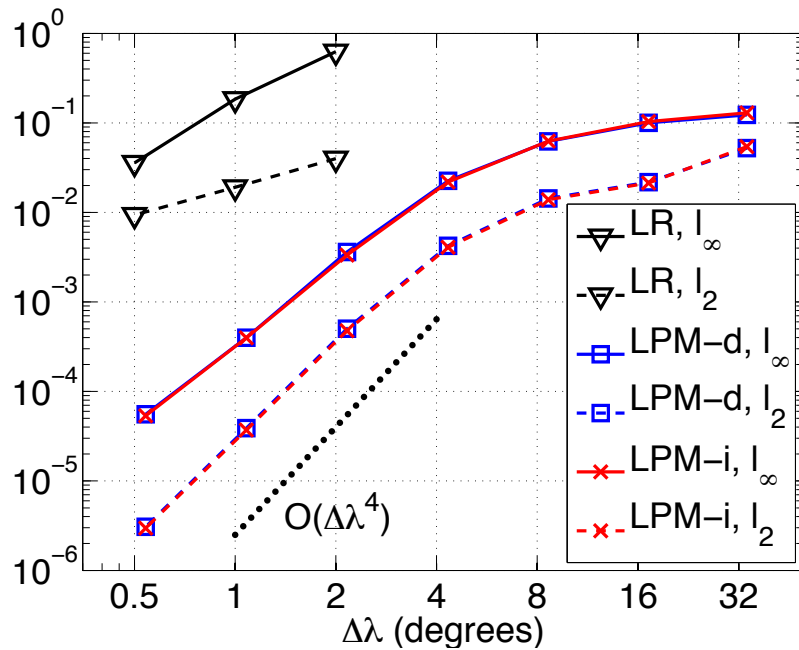
test case 1 : moving vortices flow



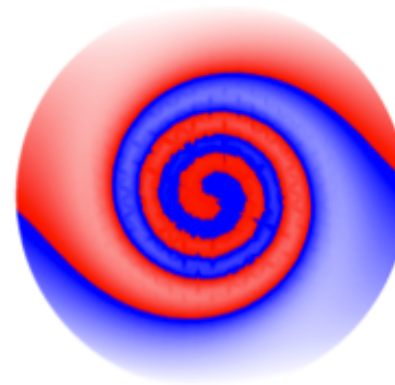
movie

tracer error

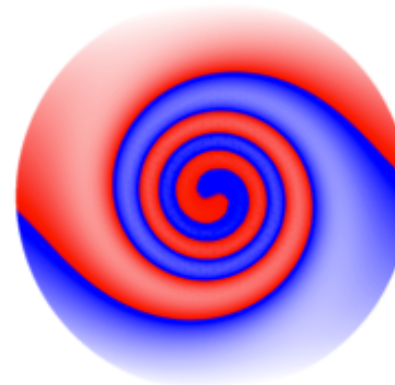
test case 1



$$\Delta\lambda = 8.64^\circ$$



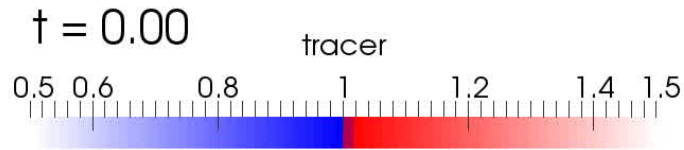
$$\Delta\lambda = 4.33^\circ$$



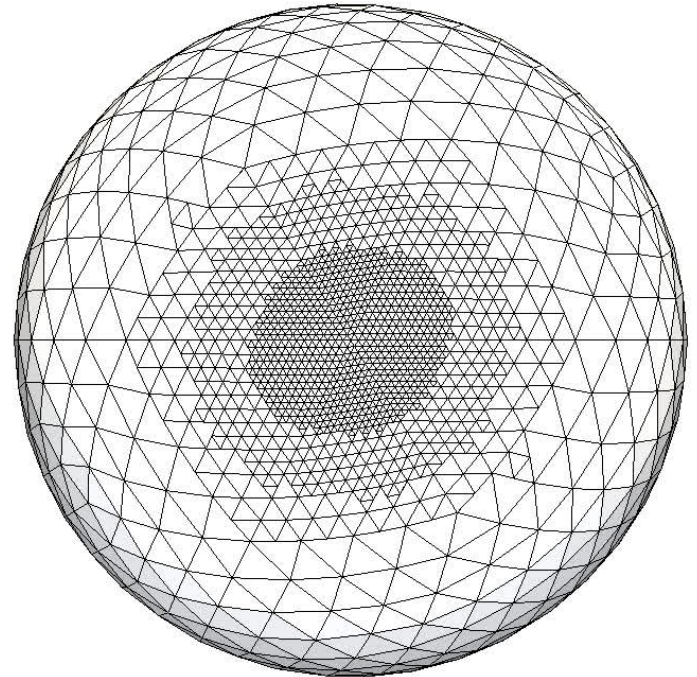
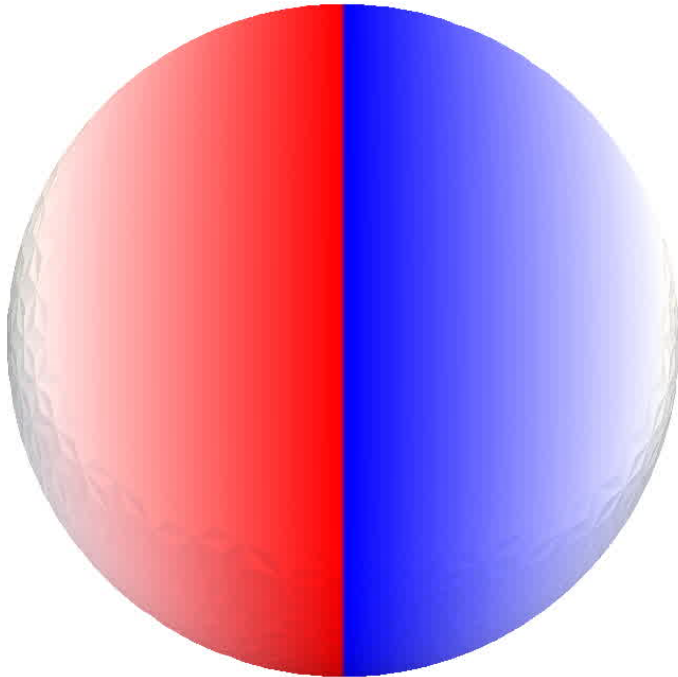
$$\Delta\lambda = 2.16^\circ$$

test case 1 : moving vortices flow

- adaptive panel refinement

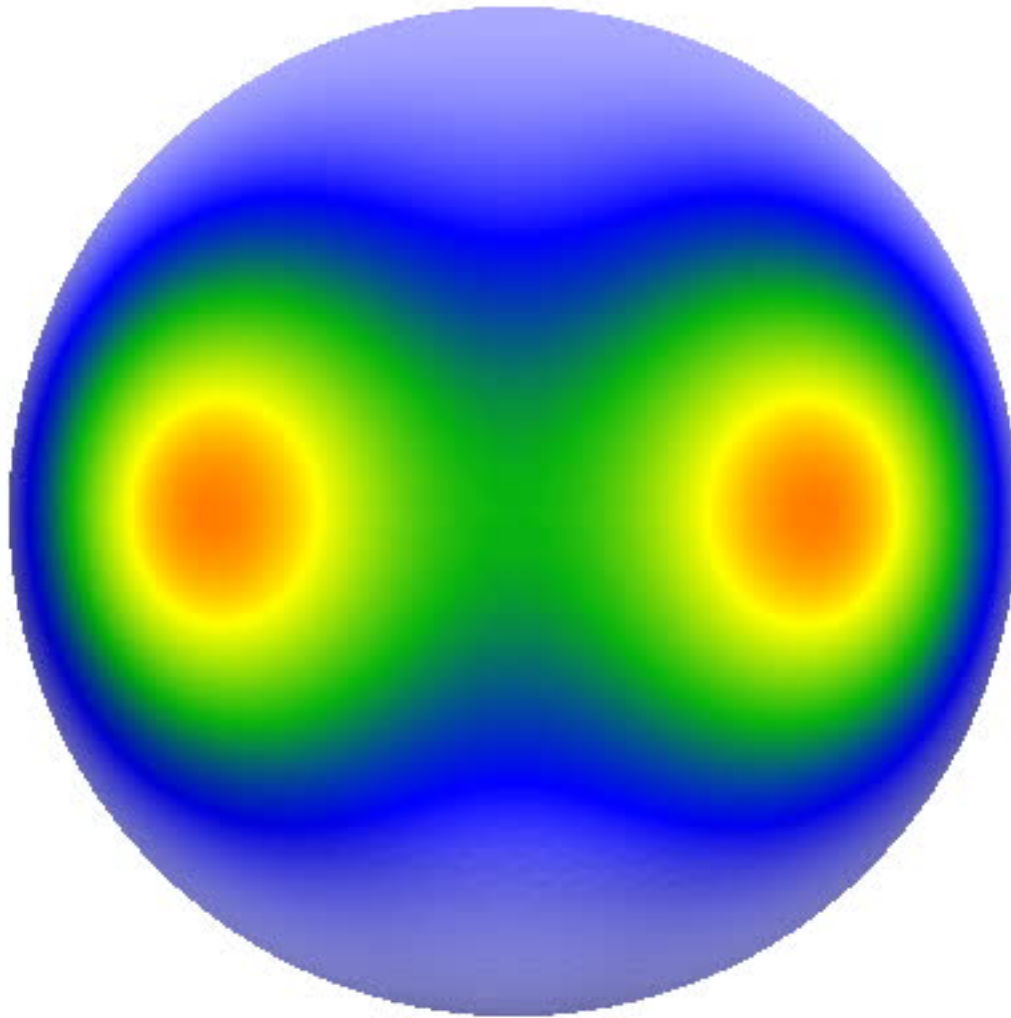


$$\zeta(\mathbf{x}_k) A_k < \epsilon \Gamma_{\max}$$



movie

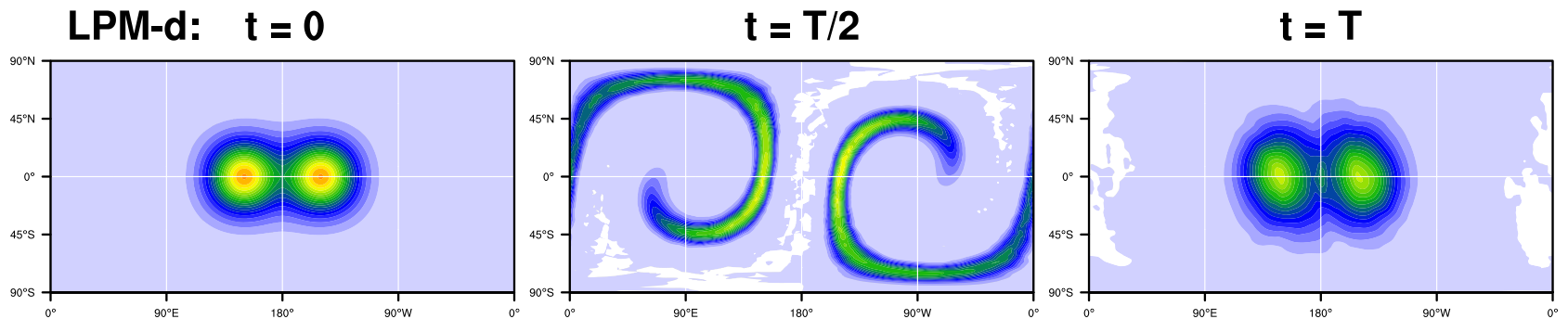
test case 2 : reversing-deformational flow
- Gaussian hills tracer



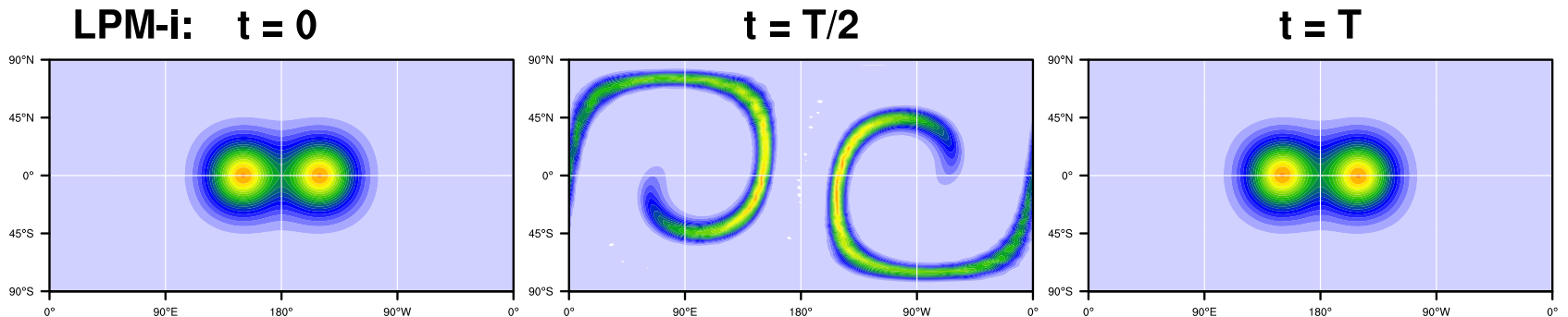
test case 2 : reversing-deformational flow

Gaussian-hills tracer , grid spacing $\Delta\lambda = 8.64^\circ$

direct remeshing

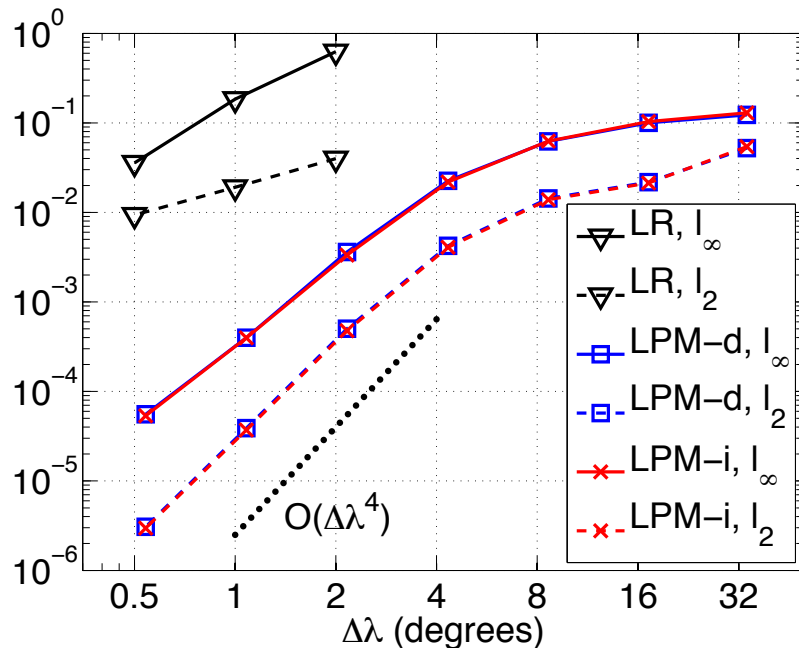


indirect remeshing

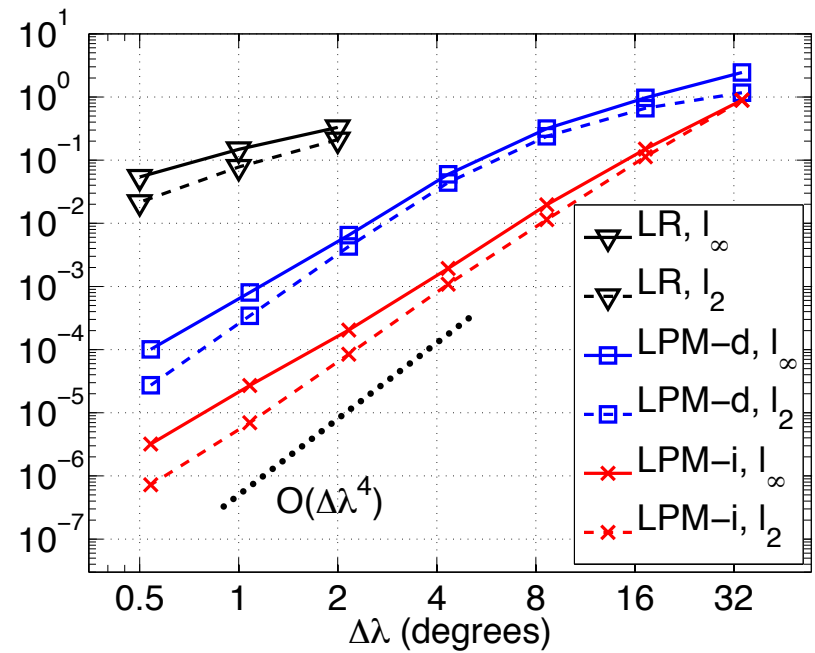


tracer error

test case 1



test case 2



vortex dynamics on a sphere (partial list!)

point vortices

Kimura-Okamoto (1987), Newton-Sakajo (2007) ...

vortex patches

Dritschel (1989, 2004), Dritschel-Polvani (1992)

Crowdy-Cloke (2003), Surana-Crowdy (2008) ...

vortex sheets

Sakajo (2009) ...

present work : smooth vorticity

barotropic vorticity equation : Eulerian form

velocity : $\mathbf{u} = \nabla\psi \times \mathbf{x}$

vorticity : $\zeta = \nabla \times \mathbf{u}$

Poisson equation : $\nabla^2\psi = -\zeta$

vortex dynamics : $\frac{\partial\zeta}{\partial t} + \mathbf{u} \cdot \nabla(\zeta + f) = 0$

Coriolis parameter : $f = 2\Omega \sin \theta$

barotropic vorticity equation : Lagrangian form

flow map : $\mathbf{x} = \mathbf{x}(\vec{\alpha}, t)$

vorticity : $\zeta = \zeta(\mathbf{x}, t)$

stream function : $\psi(\mathbf{x}) = -\frac{1}{4\pi} \int_S \log(1 - \mathbf{x} \cdot \tilde{\mathbf{x}}) \zeta(\tilde{\mathbf{x}}) dA(\tilde{\mathbf{x}})$

Biot-Savart law : $\frac{d\mathbf{x}}{dt} = -\frac{1}{4\pi} \int_S \frac{\mathbf{x} \times \tilde{\mathbf{x}}}{1 - \mathbf{x} \cdot \tilde{\mathbf{x}}} \zeta(\tilde{\mathbf{x}}, t) dA(\tilde{\mathbf{x}})$

vortex dynamics : $\frac{d\zeta}{dt} = -2\Omega \frac{dz}{dt}$

discretization

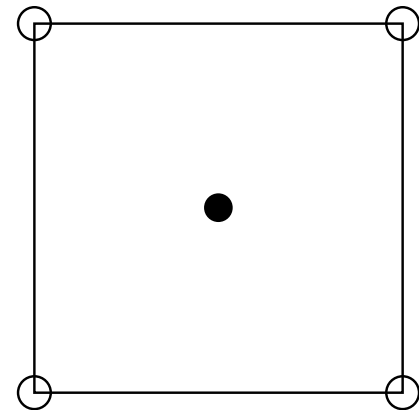
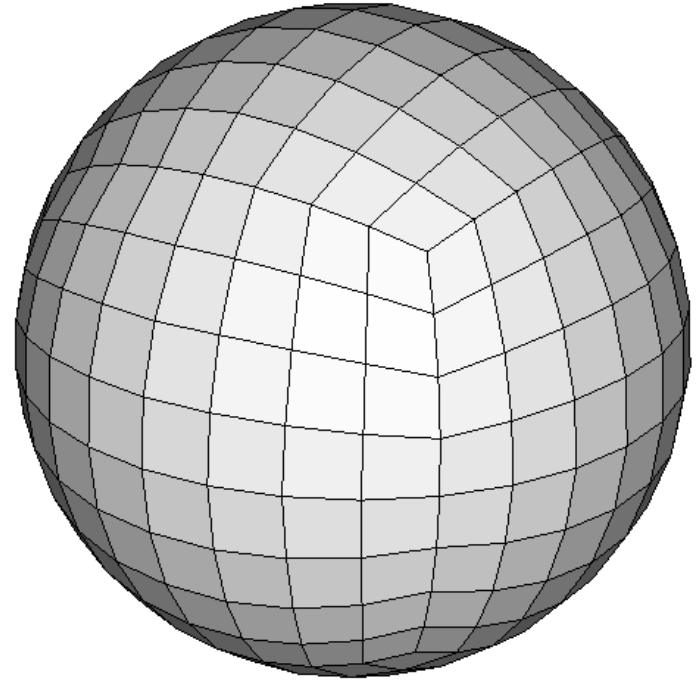
$$\mathcal{S} = \cup_{k=1}^N A_k$$

panels : cubed sphere

particles : $\mathbf{x}_j(t) = \mathbf{x}(\vec{\alpha}_j, t)$

$$\frac{d\mathbf{x}_j}{dt} = -\frac{1}{4\pi} \sum_{\substack{k=1 \\ k \neq j}} \frac{\mathbf{x}_j \times \mathbf{x}_k}{1 - \mathbf{x}_j \cdot \mathbf{x}_k} \zeta_k A_k$$

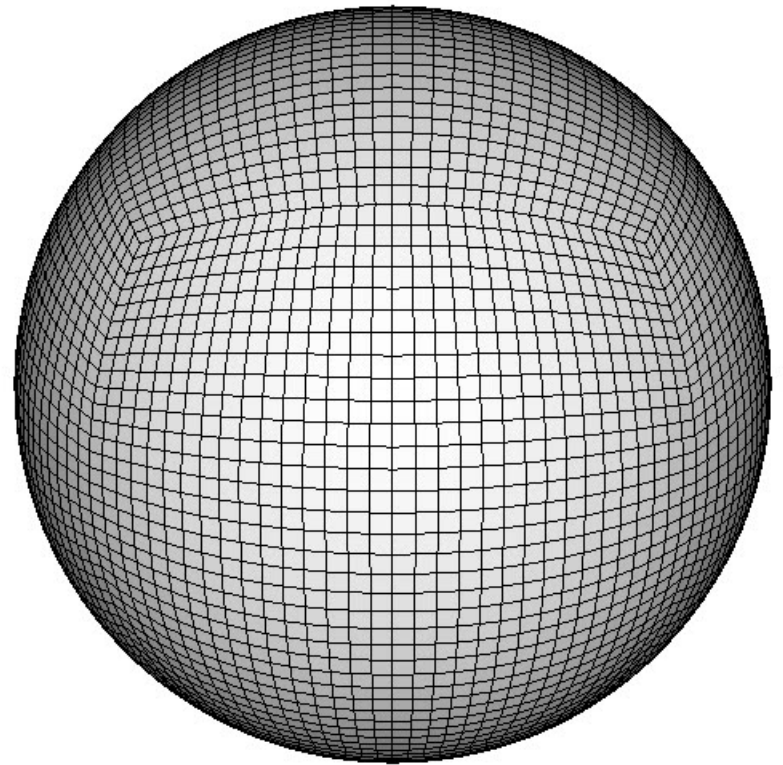
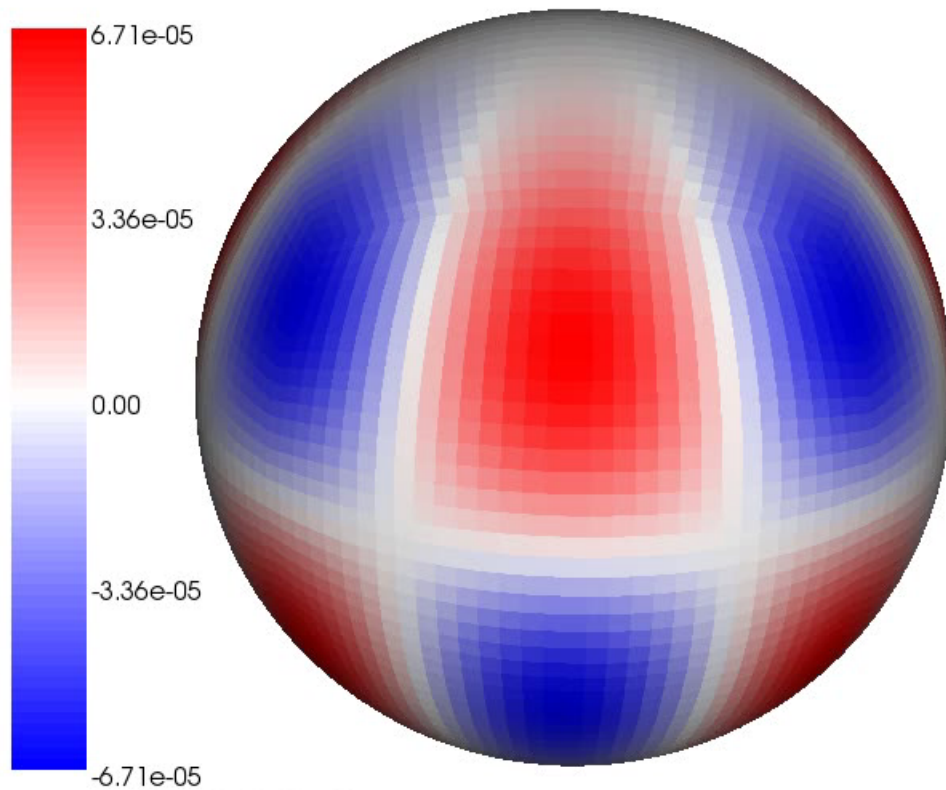
vorticity : $\frac{d\zeta_j}{dt} = -2\Omega \frac{dz_j}{dt}$



Rossby-Haurwitz wave – no remeshing

vorticity

panels



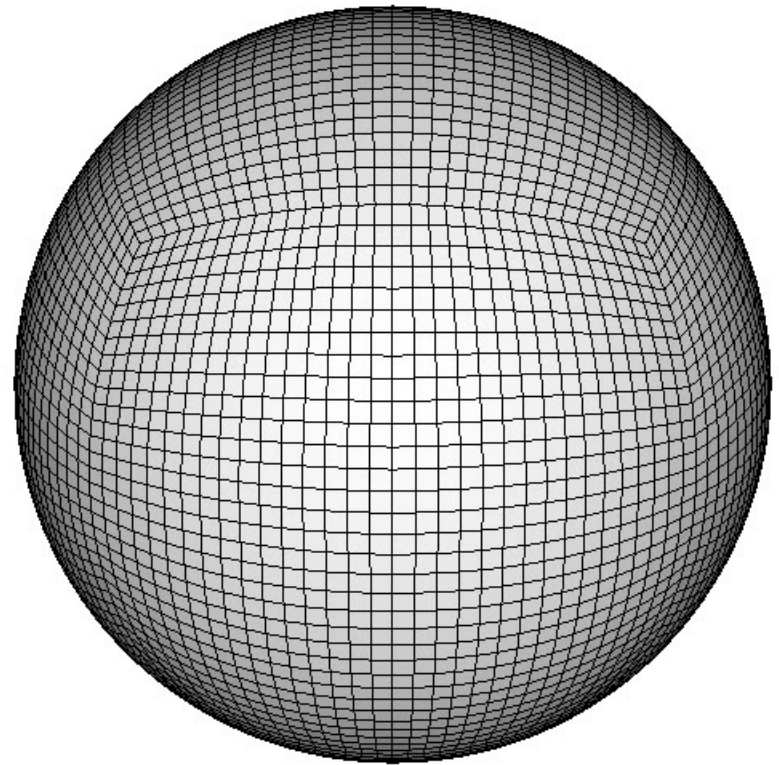
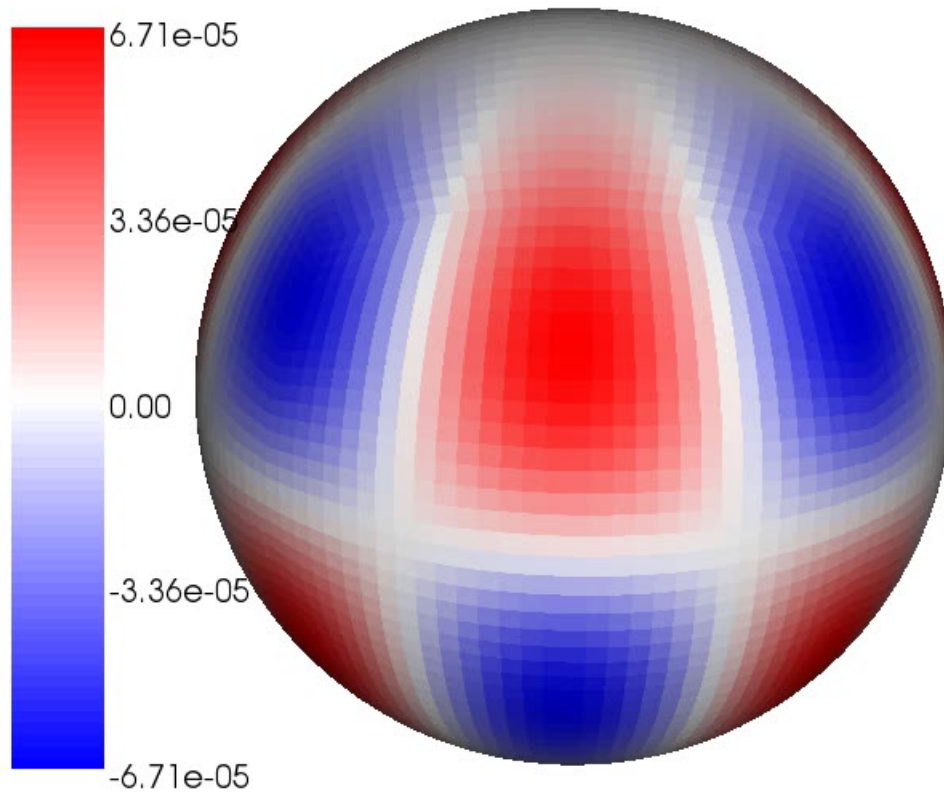
$\dagger = 0.0000, N = 6144$

movie

Rossby-Haurwitz wave – remeshing

vorticity

panels



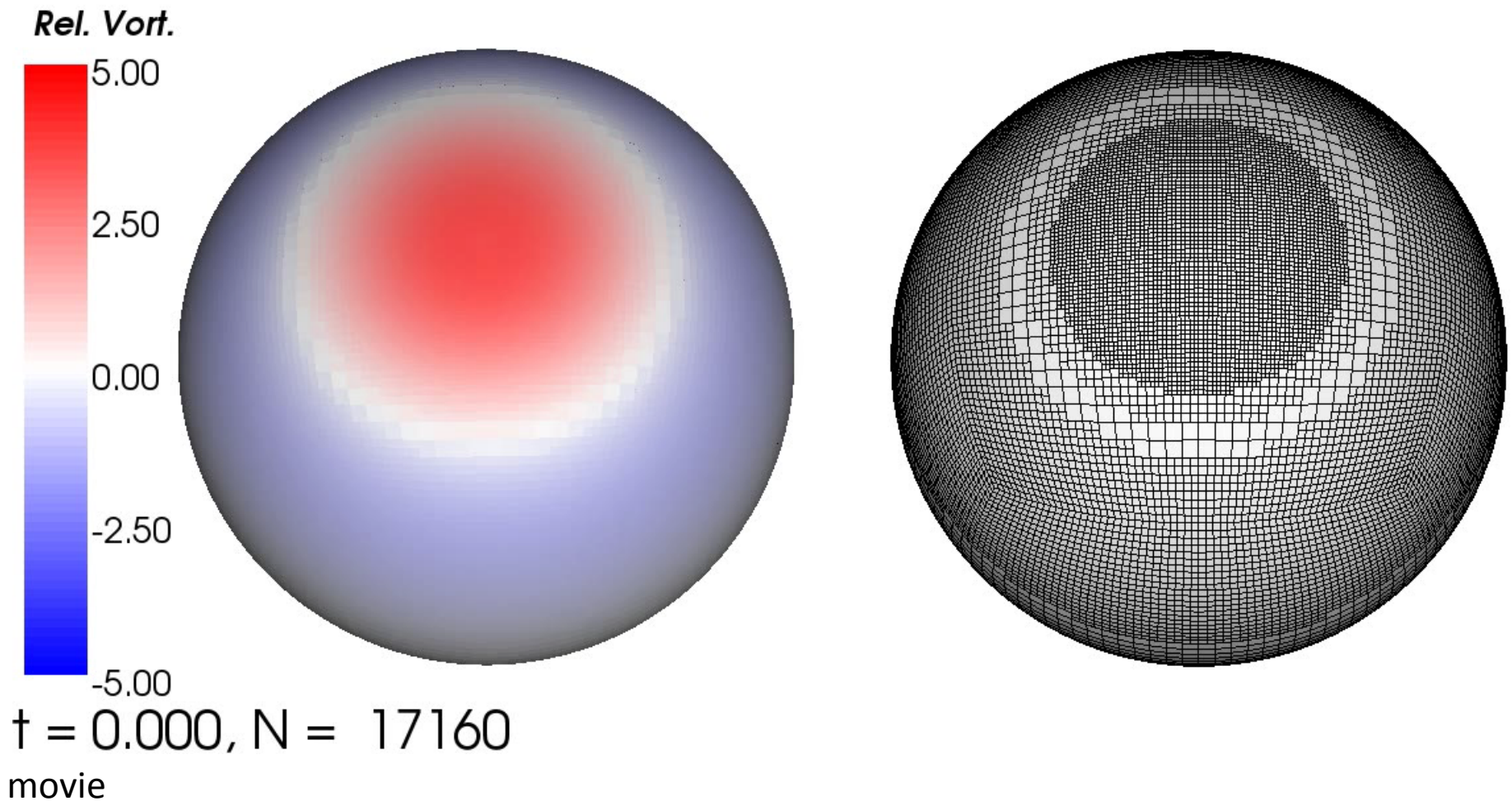
$t = 0.00, N = 6144$

movie

disruption of the polar vortex

motivated by Jukes & McIntyre (1987)

remesh + refine



fluid flow in 2D : incompressible, inviscid

Eulerian form

$$\frac{\partial \omega}{\partial t} + \mathbf{u} \cdot \nabla \omega = 0, \quad \nabla \cdot \mathbf{u} = 0$$

$$\mathbf{u} = \nabla^\perp \psi, \quad \nabla^2 \psi = -\omega$$

Lagrangian form

$$\mathbf{x}(\mathbf{a}, t), \quad \frac{d\mathbf{x}}{dt} = \mathbf{u}(\mathbf{x}, t), \quad \mathbf{x}(\mathbf{a}, 0) = \mathbf{a}$$

$$\mathbf{u}(\mathbf{x}, t) = \int K_\delta(\mathbf{x}, \mathbf{y}) \omega(\mathbf{y}, t) d\mathbf{y}, \quad K_\delta(\mathbf{x}, \mathbf{y}) = -\nabla^\perp \frac{\ln(|\mathbf{x} - \mathbf{y}|^2 + \delta^2)}{4\pi}$$

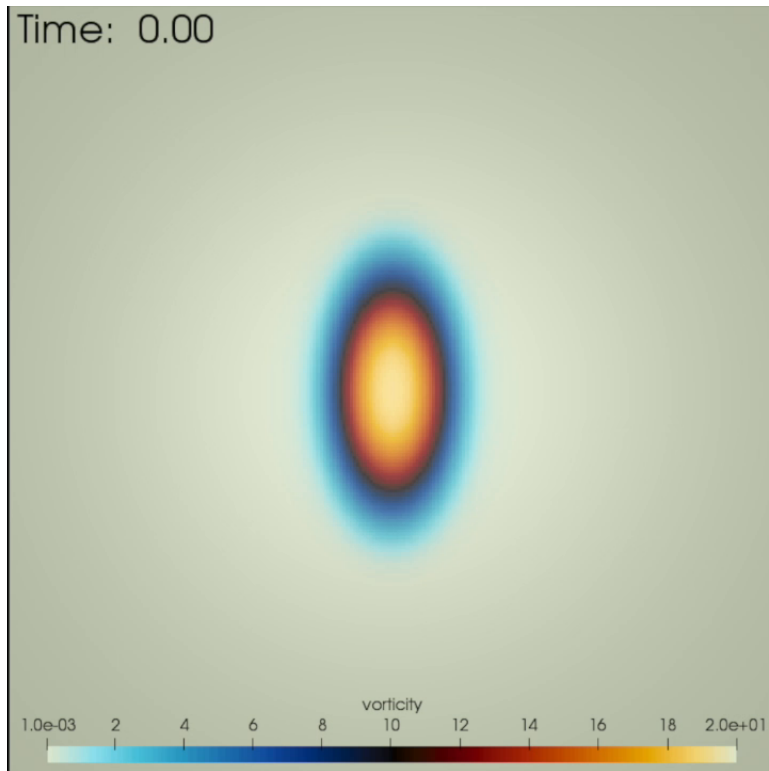
$$\omega(\mathbf{x}, t) = \omega_0(\mathbf{a})$$

- particle/panel method + remeshing, AMR, treecode

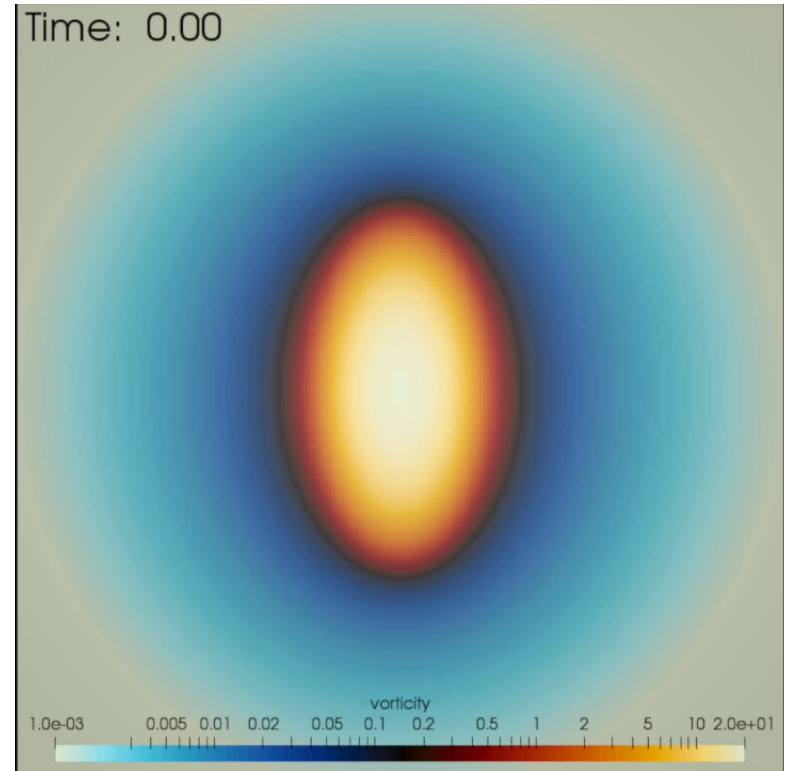
elliptic vortex

Koumoutsakos (1997)

linear scale

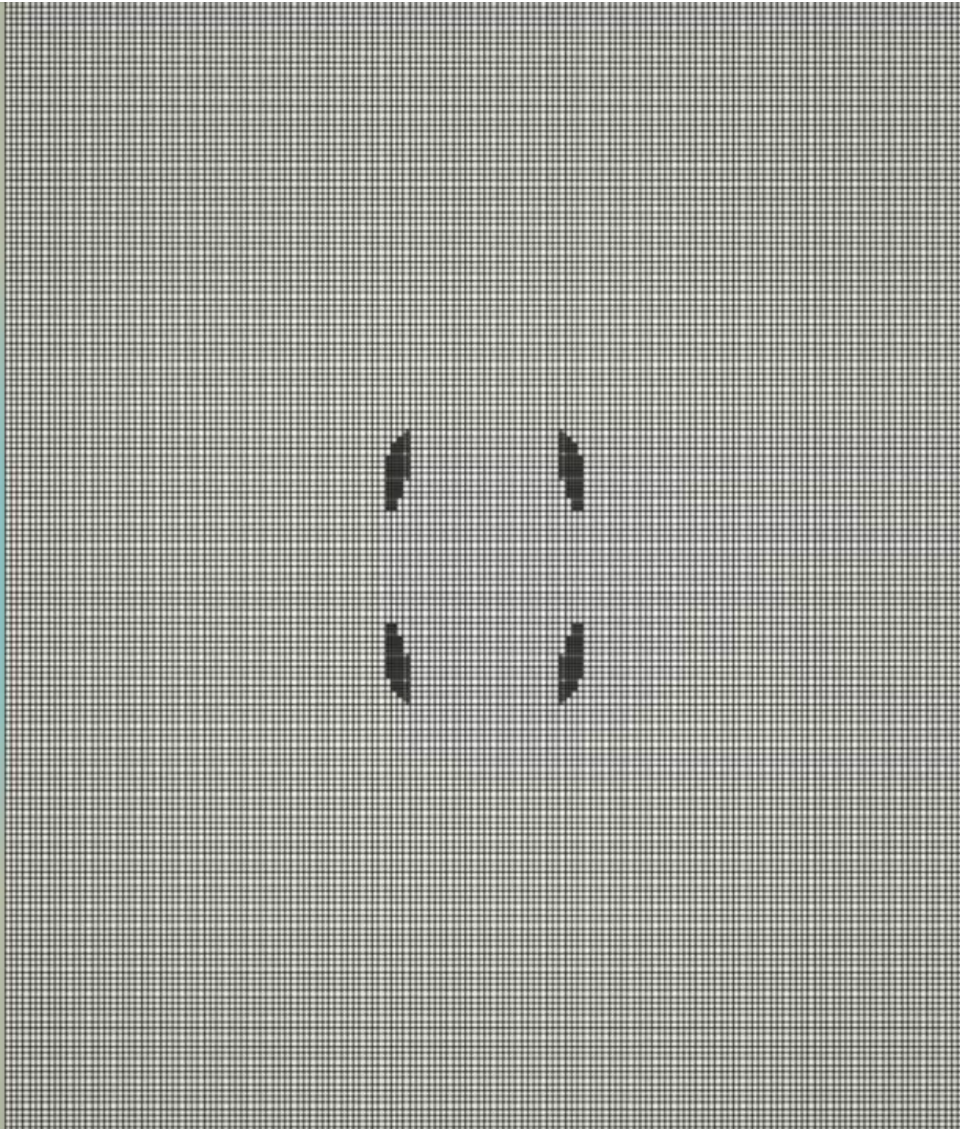
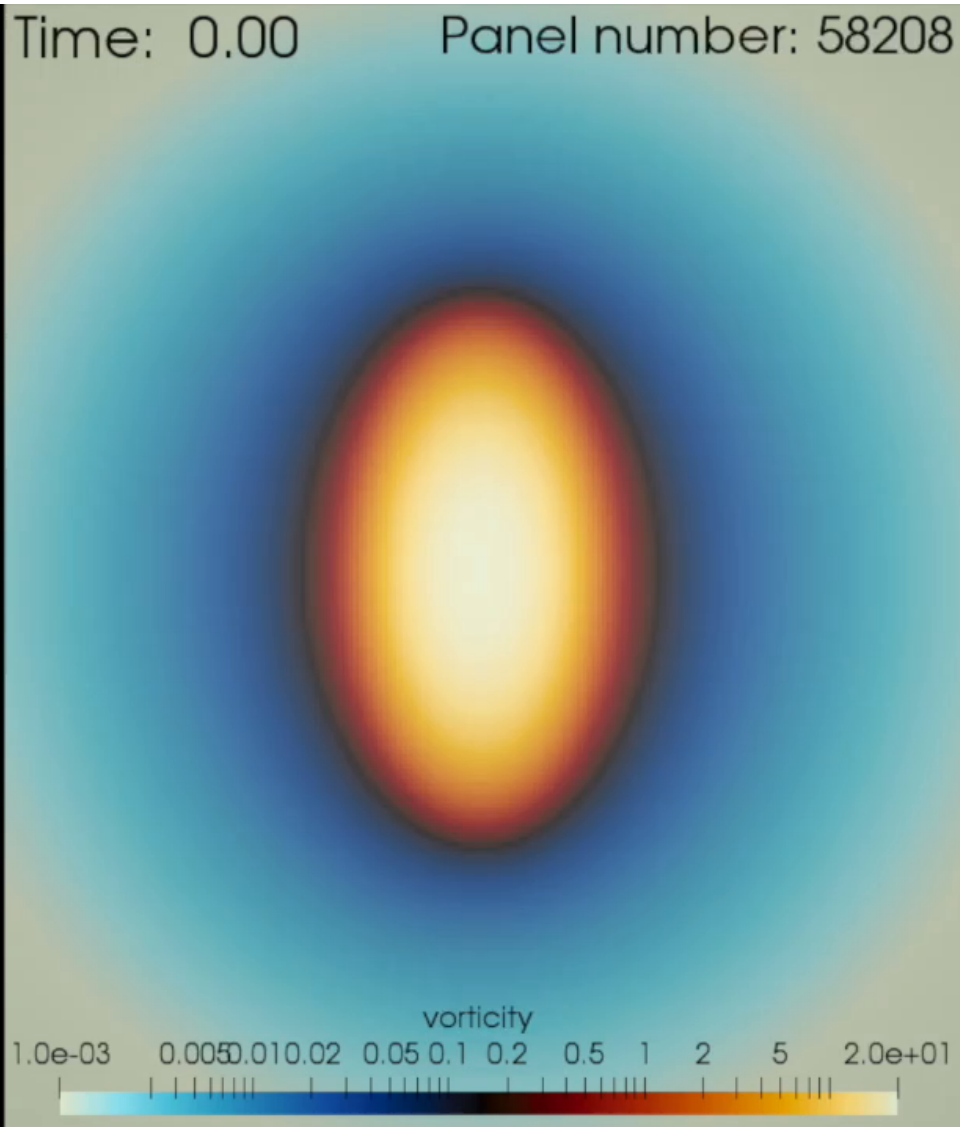


logarithmic scale



elliptic vortex

adaptive mesh

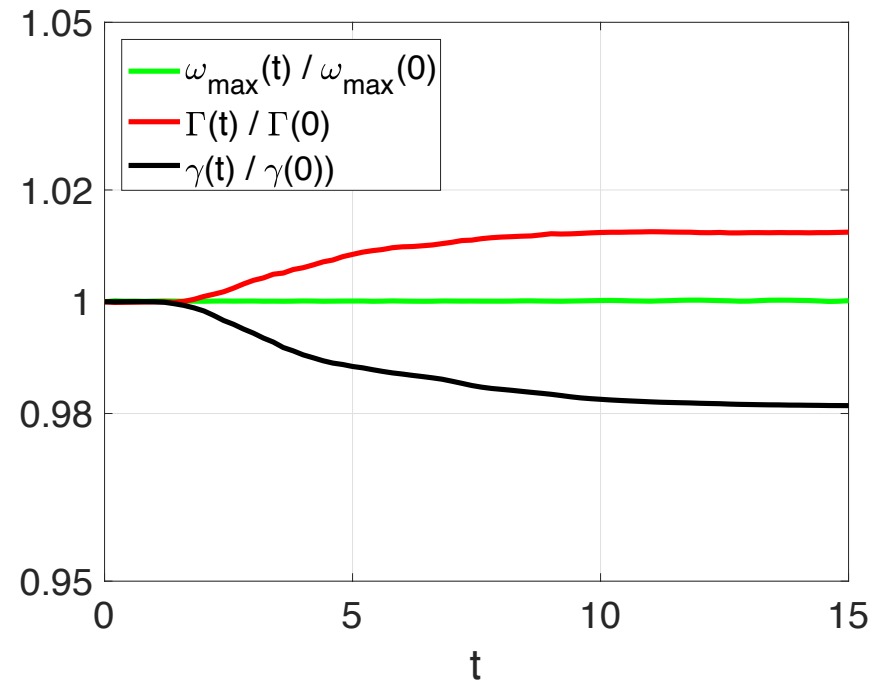
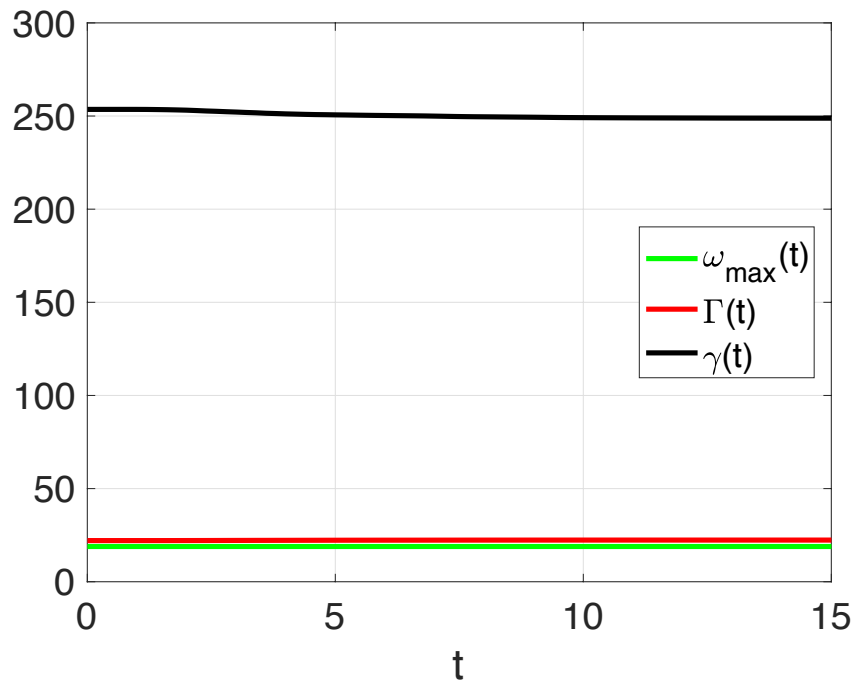


elliptic vortex 1 : conserved quantities

maximum vorticity : $\omega_{\max}(t) = \max_{\mathbf{x}} \omega(\mathbf{x}, t)$

circulation : $\Gamma(t) = \int \omega(\mathbf{x}, t) dA$

enstrophy : $\gamma(t) = \int \omega^2(\mathbf{x}, t) dA$



thanks to my collaborators -

thanks to my collaborators -

Weihua Geng



Peter Bosler



Ling Xu



current/future projects

protein-solvent electrostatics

- binding energy of protein-ligand complex
- protein pKa

incompressible fluid flow

- viscosity, 3D, solid boundaries
- shallow water equations